## **<https://doi.org/10.5281/zenodo.10395829>**

## **Alimov Z Turg'unov M**

**Abstract**:*In this paper, we construct a mathematical model of a damped using two different wave equation. The process is linked with vibration of stretched string taking place in a special medium, where damping force acting in a short time and then disappears for the rest of the time. The key tool is a method of a separation of variables.*

**Keywords:** *wave equation, damping force, method of separation variables.*

It is well-known that motion caused by vibration of string due to stretch can be modeled by wave equation  $\lceil 1 \rceil$ . If it happens in a special media, where damping force is acting then model wave equation will be modified adding  $\frac{\partial u}{\partial x}$ *t*  $\hat{o}$  $\partial$ term. There are many works devoted to the studying spring-mass-damping processes, related with integer or fractional order differential equations [2-4].

Let us consider the following problem:

Problem. A string is stretched and secured on the x-axis at  $x = 0 = 0$  and  $x = \pi$  for some time-interval, for instance,  $t \in (0,T)$  , where  $T>0$  is a real number.

If the transverse vibrations take place in a medium that imports a resistance proportional to the instantaneous velocity depending on time and the initial displacement is  $\varphi$ (x) find the equation of the motion.

 A mathematical model of this process is based on wave equation of the following form

$$
\frac{\partial^2 u}{\partial t^2} + \beta(t) \frac{\partial u}{\partial t} - \frac{\partial^2 u}{\partial x^2} = 0,
$$
\n(1)

where  $\beta(t)$  will be defined as

$$
\beta(t) = \begin{cases} A, & 0 < t < T_0 \\ 0, & T_0 < t < T, \end{cases}
$$

*A* is a real constant. Boundary conditions will have a form of

$$
u(t,0) = 0, \ \ u(t,\pi) = 0, \ \ 0 \le t \le T. \tag{2}
$$

Here we will impose initial conditions

$$
u(0,x) = \varphi(x), \ 0 \le x \le \pi, \ u_t(0,x) = 0, 0 < x < \pi. \tag{3}
$$

The domain of consideration is a rectangle  $\Omega = \{(t, x): 0 < x < \pi, 0 < t < T\}$ . Since  $\beta(t)$  is changing its value depending on  $t$ , we need to divide  $\Omega$  into two parts, namely, Volume. 6, Issue 08, December (2023)

blume. 6, Issue 08, December (2023)<br>and  $\Omega_2 = \{(t, x): 0 < x < \pi, T_0 < t < T\}$ . Therefore, equation (1)

reads as

as  
\n
$$
\frac{\partial^2 u}{\partial t^2} + A \frac{\partial u}{\partial t} - \frac{\partial^2 u}{\partial x^2} = 0, \quad 0 < t < T_0,
$$
\n
$$
\frac{\partial^2 u}{\partial t^2} - \frac{\partial^2 u}{\partial x^2} = 0, \quad T_0 < t < T.
$$
\n(4)

On the line 
$$
t=T_0
$$
 one needs to consider conjugating conditions  

$$
u(T_0+0,x) = u(T_0+0,x), 0 \le x \le \pi,
$$
 (5)

$$
\left.\frac{\partial u}{\partial t}\right|_{t=T_0-0} = \frac{\partial u}{\partial t}\Big|_{t=T_0+0}, 0 < x < \pi.
$$
\n(6)

Our aim is to find a function  $u(t, x)$ , which satisfies equation (4) together with boundary conditions (2), initial conditions (3) and conjugation conditions (5)-(6). This

function should also satisfy the following regularity conditions.  
\n
$$
u(t, x) \in C(\overline{\Omega}) \cap C^2(\Omega_1 \cup \Omega_2) \cap C^1(\Omega)
$$
. (7)

We will search for nontrivial solution  $u(t, x)$  as follows:

$$
u(t,x) = u(t) \cdot X(x). \tag{8}
$$

Substituting (8) into (4) and (2) we will get the following eigenvalue problem in *x*variable:

$$
\begin{cases} X''(x) + \lambda X(x) = 0 \\ X(0) = X(T) = 0, \end{cases}
$$
 (9)

where  $\lambda$  is a real number.

It is well-known that the problem (9) has nontrivial solutions at 
$$
\lambda = n^2
$$
, n=1,2,... [1]  
 $X_n(x) = \sin nx$ . (10)

In time-variable we will get the following equations:

$$
u'' + Au' + \lambda u = 0, 0 < t < T_0,
$$
  
\n
$$
u'' + \lambda u = 0, T_0 < t < T.
$$
\n(12)

General solution of (11) and (12) have a form  
\n
$$
u(t) = e^{-\frac{A}{2}t} \left( C_1 \cos \frac{\sqrt{4\lambda - A^2}}{2} t + C_2 \sin \frac{\sqrt{4\lambda - A^2}}{2} t \right), \qquad -2\sqrt{\lambda} < A < 2\sqrt{\lambda},
$$

$$
u(t) = C_3 \cos \sqrt{\lambda} t + C_4 \sin \sqrt{\lambda} t.
$$

63 1 ( , ) : 0 , 0 *t x x t T* 0 Considering that 2 *n* , we have 2 2 4 4 <sup>2</sup> ( ) C cos sin , 1 2 2 2 *n n n At A A u t e t C t* 2 2, *A* (13) 3 4 ( ) cos sin *n n n u t C nt C nt* . (14)

According to the superposition principle, we have

Volume. 6, Issue 08, December (2023)

$$
u(t,x) = \sum_{n=1}^{\infty} u_n(t) \sin nx,
$$
 (15)

where  $u_n(t)$  is defined by (13) for  $0 < t < T_0$  and by (14) for  $T_0 < t < T$ .

In order to find unknown constants  $C_{in}$   $(i=1,4)$ , we will use initial conditions (3) and conjugating conditions (5)-(6).

After the finding coefficients  $C_{_{in}}$ , we need to prove the uniform convergence of series (15) and other infinite series corresponding to functions  $u_{_n}\big(t,x\big),u_{_{\mathrm{xx}}}\big(t,x\big).$  During this process we have to impose certain conditions to given data in order to guarantee that found function  $u(t, x)$  will satisfy regularity conditions (7).

## **REFERENCES:**

1. A.N. Tikhonov, A.A. Samarskii. Equation of Mathematical Physics. Courier Corporation, 2013, 800 p.

2. J. F. Gomez-Aguilar, H. Yepez-Martinez, C. Calderon-Ramon, I. Cruz-Ordunia, R. F. Escobar-Jimenez and V. H. Olivares-Peregrino, Modeling of a mass-spring-damper system by fractional derivatives with and without a singular kernel, Entropy, 17 (2015), pp.6289- 6303.

3. J. F. Gomez-Aquilar, J. J. Rosales-Garcia, J. J. Bernal-Alvarado, T. Cordova-Fraga and R. Guzman-Cabrera, Fractional mechanical oscillators, Rev. Mexicana Fisica, 58(2012), pp.348-352.

4. N. Al-Salti, E.Karimov, K.Sadarangani. On a differential equation with Caputo-Fabrizio fractional derivative of order  $1 < \beta < 2$  and application to mass-spring-damper system.